

On a possibility to combine the order effect with sequential reproducibility for quantum measurements

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Abstract

In this paper we study the problem of a possibility to use quantum observables to describe a possible combination of the order effect with sequential reproducibility for quantum measurements. By the order effect we mean a dependence of probability distributions (of measurement results) on the order of measurements. We consider two types of the sequential reproducibility: adjacent reproducibility ($A - A$) and separated reproducibility ($A - B - A$). The first one is reproducibility with probability 1 of a result of measurement of some observable A measured twice, one A measurement after the other. The second one, $A - B - A$, is reproducibility with probability 1 of a result of A measurement when another quantum observable B is measured between two A 's. Heuristically, it is clear that the second type of reproducibility is complementary to the order effect. We show that, surprisingly, for an important class of quantum observables given by positive operator valued measures (POVMs), this may not be the case. The order effect can coexist with a separated reproducibility as well as adjacent reproducibility for both observables A and B . However, the additional constraint in the form of separated reproducibility of the $B - A - B$ type makes this coexistence impossible. Mathematically, this paper

is about the operator algebra for effects composing POVMs. The problem under consideration was motivated by attempts to apply the quantum formalism outside of physics, especially, in cognitive psychology and psychophysics. However, it is also important for foundations of quantum physics as a part of the problem about the structure of sequential quantum measurements.

1 Introduction

Sequential measurements play an important role in quantum measurement theory, e.g., [1]–[5]. In particular, sequential measurements of two incompatible observables A and B induce the probability order effect which gives the probabilistic representation of the principle of complementarity, in the form of nonexistence of the joint probability distribution of such observables. In other words, sequential probability distributions $p_{A-B}(x, y)$ and $p_{B-A}(y, x)$ do not coincide.

We remark that if A and B are quantum measurements of the von Neumann-Lüders type (measurements of the first kind), i.e., they are mathematically represented by Hermitian operators and the state transformers (quantum operations)¹ are given by the orthogonal projectors corresponding to the eigenvalues², then they have the property of *adjacent sequential reproducibility*: for $A - A$ and $B - B$ measurements the values observed in the first measurement are reproduced (with probability 1) in the second measurement. Thus, for quantum measurements of the first kind, the condition of adjacent reproducibility is a redundant constraint. However, for quantum measurements of the second kind, i.e., those which cannot be represented as measurements of the first kind, this constraint is nontrivial.

We remark that, for quantum measurements of the second kind neither an observable nor a state transformer have to be represented by orthogonal projectors. In general, for measurements of the second kind, observables are represented by POVMs. We also remark that a quantum measurement with

¹The terminology “state transformer” [2] is may be old fashioned. In modern literature [7] the terms of quantum operation and quantum channel are typically in the use. The latter notions are motivated by quantum information theory [7]. However, the problems handled in this paper have no direct relation to quantum information theory. They are motivated by applications of the quantum calculus of probabilities outside of physics, see the last part of this section for details. In such applications the terminology “state transformer” seems to be more appropriate.

²In the mathematical model, see, e.g., [1]–[5], quantum measurement is represented with the aid of two structures: a quantum observable (Hermitian operator or more generally POVM) and a state transformer.

an observable given by a Hermitian operator can also be of the second kind: if the corresponding state transformers are not given by orthogonal projectors. The latter class of quantum measurements will play an important role in problems under the study in this paper.

For measurements of the first kind, the order effect can be approached only in case of incompatibility: commutativity of the observables ($[A, B] = 0$) implies the existence of the joint probability distribution serving both sequential measurements, $A - B$ and $B - A$. Intuitively, in the case of incompatible observables a measurement of the observable B after a preceding measurement of the observable A (with the result $A = x$) modifies crucially the post-measurement state ψ_A generated as the result of the A measurement. Thus, in the post-measurement state after the sequence of measurements $A - B$, the information about the value $A = x$ is at least partially washed out and one cannot expect that in the sequence of measurements $A - B - A$ the value $A = x$ would be reproduced (obtained in the second A measurement) with probability 1. Thus, one would expect that the quantum order effect cannot be combined with both the adjacent and separated reproducibility. In paper [8] it was shown that, for measurements of the first kind, this is really the case.

The problem of an extension of this result to measurements of the second kind is mathematically more complicated. Some special class of measurements of the second kind was considered in the aforementioned paper and it was shown that the combination of $A - A$, $B - B$, and $A - B - A$ reproducibilities implies disappearance of the order effect. The class of measurements of the second kind considered in [8] is characterized by the following conditions:

C0 Observables are represented by POVMs.

C1 The state transformer corresponding to an effect E has the form

$$\psi \rightarrow \frac{M\psi}{\|M\psi\|}, \quad (1)$$

where

$$E = M^*M \quad (2)$$

is some representation of the effect.

C2 If $E = P$ is an orthogonal projector, then its state transformer is given by this projector

$$\psi \rightarrow \frac{P\psi}{\|P\psi\|}. \quad (3)$$

The condition **C1**, although it restricts the class of possible state transformers, is very natural, because the majority of state transformers used in

applications really have this simple form, (1). By the condition **C2** the class of measurements corresponding to the effects given by projectors coincides with the class of measurements of the von Neumann-Lüders type (measurements of the first kind). Although this condition is also quite natural, it restricts essentially the class of state transformers. In this paper we want to proceed with the class of quantum measurements described solely by the conditions **C0, C1**. Thus, even for an effect given by an orthogonal projector, in general the corresponding state transformer does not coincide with this projector.

Surprisingly, this extension of the class of quantum measurements makes the situation essentially more complicated (and interesting). The order effect can coexist with separated reproducibility of $A - B - A$ type as well as adjacent reproducibility for both observables A and B . However, the additional constraint in the form of separated reproducibility of the $B - A - B$ type makes this coexistence impossible.

Finally, we remark that this study (as well as the preceding study [8]) was motivated by applications of the quantum formalism outside of physics, namely, in cognitive psychology and psychophysics, see, e.g., [9]-[19] for introduction. In such applications the mathematical formalism of quantum mechanics is treated as an operational formalism, see, e.g., M. D' Ariano [20], for handling probabilistic data collected in aforementioned domains of research. This formalism works well for a wide class of psycho-effects, e.g., the disjunction effect and the order effect. In particular, to represent these two effects one can use the von Neumann-Lüders measurements. As is well known, such observables satisfy the condition of adjacent reproducibility, i.e., $A - A$ reproducibility, which is natural for the majority of experiments of cognitive psychology. However, as was pointed out in [8], in psychophysics one can find the experimental situations violating adjacent reproducibility. Therefore one has to proceed with generalized observables given by POVMs.³ And this was the natural step in development of quantum-like modeling in humanities.

As was first understood by the authors of [8], there was a hidden pitfall in the rapidly increasing stream of applications of the quantum formalism to humanities, namely, the problem of separated reproducibility. As was pointed out in [8], some experimental contexts of cognitive psychology are characterized by separated reproducibility effect in combination with the

³Observables of this class appeared in the applications of the quantum mathematical formalism to humanities even by another reason [10]: in experiments performed, e.g., in cognitive psychology, the matrices of transition probabilities for observables with non-degenerate spectra are typically not double-stochastic. However, for the von Neumann-Lüders observables, they should be double-stochastic.

order effect. For example, political opinion polls often demonstrate order effects, but here, e.g., by replying “yes” to the first question A , and “no” to the second question B a respondent is typically “firm in her preferences” expressed in the form of the A -yes, so she will practically definitely say “yes” again if asked the A -question again, i.e., in the $A - B - A$ experiment. We also remark that in known and thinkable experimental contexts in cognitive psychology, the same should also happen for the $B - A - B$ experiment. As was shown in [8], it is impossible to use the measurements of the first kind to describe such a situation. The following natural question was posed by the authors of [8]: *Is the operational quantum formalism powerful enough to cover all possible experimental contexts arising in humanities?* In this paper, we continue to work to find an answer to this question. And for the moment, the answer is negative. However, although this paper covers a wide and natural class of state transformers, it is still not the most general. There is still a “loophole in the proof” that the power of quantum methods in humanities is restricted. We remark that in principle there is no reason to expect that the quantum-like operational formalism serving humanities would coincide with the quantum physical formalism. It may happen that novel quantum-like models would be explored, cf. [10].

Although the problem of combination of the order effect with the two types of reproducibility was motivated by applications of the quantum formalism outside of physics, mainly cognitive psychology, it is also important for quantum foundations as a part of the problem about the structure of sequential quantum measurements, cf. [5]. There is also an important experimental dimension: *Can one find physical measurements exhibiting combination of the order effect with adjacent and separated reproducibilities?*

In this paper we consider only the case of finite dimensional state spaces. (Only such state spaces are used up to now in applications to humanities.) The situation in the infinite-dimensional case is very different, see, e.g., Proposition 8, and this case has to be studied separately.

2 The basic consequence of adjacent reproducibility

In this section we show that the class of effects corresponding to measurements determined by the conditions **C0**, **C1**, and satisfying the condition of adjacent reproducibility coincides with the class of orthogonal projectors. Hence, the class the corresponding POVMs coincides with the class of von Neumann-Lüders observables. But, in general, the state transformers are

not reduced to projectors. Here the basic result is presented in Theorem 1. (This statement was formulated in [8]. However, its proof contained a loophole which could not be closed.)

Theorem 1. *Let E be a Hermitian operator such that $0 \leq E \leq I$ (an effect) and let $\langle E\phi, \phi \rangle = 1$ for some pure state ϕ . Then this ϕ is an eigenvector of E with the eigenvalue $\lambda = 1$.*

Proof. Consider in H a basis (e_j) consisting of eigenvectors of E , i.e., $Ee_j = \lambda_j e_j$. Here $0 \leq \lambda_j \leq 1$. We have $E\phi = \sum_j \lambda_j \phi_j e_j$, where $\phi_j = \langle \phi, e_j \rangle$. We have $\langle \phi, \phi \rangle = \sum_j |\phi_j|^2 = 1$. We also have $\langle E\phi, \phi \rangle = \sum_j \lambda_j |\phi_j|^2 = 1$. Set $O_\phi = \{j : \phi_j \neq 0\}$ (this set depends on E). Then we have the system of two quadratic equalities:

$$\sum_{j \in O_\phi} |\phi_j|^2 = 1, \quad (4)$$

$$\sum_{j \in O_\phi} \lambda_j |\phi_j|^2 = 1. \quad (5)$$

If at least one λ_j in (5) is strictly less than 1, then we come to contradiction with (5). Thus, in (5) all $\lambda_j = 1$, i.e., if $j \in O_\phi$, then $\lambda_j = 1$. Hence $E\phi = \sum_{j \in O_\phi} \phi_j e_j = \phi$.

Corollary 1. *For an effect E and the corresponding state transformer of the type **C1**, the condition of adjacent reproducibility is equivalent to the operator equality:*

$$EM = M. \quad (6)$$

Theorem 2. *For an effect E and the corresponding state transformer of the type **C1**, the operator equality (6) is equivalent to the condition “ E is an orthogonal projector”.*

Proof. Consider the representation of all operators by matrices in the orthonormal basis $(e_j, j = 1, \dots, n)$ consisting of eigenvectors of the operator $E : E = (E_{ij}) = \text{diag}(\lambda_1, \dots, \lambda_n)$, $M = (m_{ij})$. Suppose that this operator has an eigenvalue, e.g., $\lambda = \lambda_1$, which is different from 0 and 1. Consider the vector given by the first row in M , $w = (m_{11}, \dots, m_{1n})$. Then the equality (6) implies that $\lambda_1 w = w$, and, hence, $w = 0$. Thus the first row of M consists of zeros (hence, the first column of M^* also consists of zeros).

Consider the vectors $v_1 = (m_{21}, \dots, m_{n1}), \dots, v_n = (m_{2n}, \dots, m_{nn})$, the columns of M without the first zero element. Since $E = M^*M$, its matrix elements can be represented as

$$E_{ij} = \langle v_i, v_j \rangle. \quad (7)$$

Since the off-diagonal elements in E are zeros, we have: $v_1 \perp v_j, j \neq 1, \dots, v_n \perp v_j, j \neq n$. Consider the first condition of orthogonality. Here a vector in the $n - 1$ dimensional space, namely, v_1 , is orthogonal to $n - 1$ vectors. There are two possibilities: either $v_1 = 0$ or vectors v_2, \dots, v_n are linearly dependent. In the first case (7) implies that $\lambda_1 = \langle v_1, v_1 \rangle = 0$. Suppose now that $v_1 \neq 0$, and, e.g., for $k = 2$, $v_2 = \sum_{j \neq 1,2} c_j v_j$. But $v_2 \perp v_j, j \neq 2$, hence $v_2 = 0$. This means that vectors v_3, \dots, v_n are linearly dependent, e.g., $v_3 = \sum_{j \neq 1,2,3} c_j v_j$. In this way we get that v_3 is also zero and so on. Finally, we get that all vectors v_2, \dots, v_{n-1} equal to zero. The vectors v_1 and v_2 have to be orthogonal.

Now let us take again the relation $EM = M$ into account. Consider vectors $w_j = (m_{j1}, \dots, m_{jn}), j = 1, \dots, n$ (the rows of the matrix M); so the vector w used above coincides with w_1 . We have $\lambda_2 w_2 = w_2$, but $\lambda_2 = \langle v_2, v_2 \rangle = 0$, so $w_2 = 0$. In the same way we get that $w_3, \dots, w_{n-1} = 0$. Thus in the vectors v_1 and v_2 only the last coordinate can be nonzero, but such vectors can be orthogonal only if one of them is zero. The worst case would be $v_1 \neq 0$, but $v_n = 0$. But then $\lambda_n = 0$ and hence $w_n = 0$, so the last coordinate in v_1 also has to be zero.

3 The structure of state transformers for effects of the projection type

In the previous section we showed that, for measurements satisfying the condition of adjacent reproducibility, the effects composing observables, POVMs, are, in fact, orthogonal projectors. Now we plan to describe the structure of the corresponding state transformer operators, see **C1**. We shall show that they are simply the compositions of the projectors-effects with unitary operators. This mathematical fact will play the crucial role in our further studies.

We start with two simple lemmas about unitary operators which will be useful in the further considerations

Lemma 1. *Let $U : H \rightarrow H$ be a unitary operator and let X be its invariant subspace, i.e., $UX \subset X$. Then the orthogonal complement of X is also invariant subspace of U , i.e., $UX^\perp \subset X^\perp$.*

Proof. Since the kernel of a unitary operator is zero, i.e., $Ux = 0$ iff $x = 0$, dimensions of the spaces X and UX are equal. Thus, for invariant subspace, we have that $UX = X$. Hence, any $x \in X$, can be represented as $x = Ux_0, x_0 \in X$. Take any $y \in X^\perp$. Then $\langle Uy, x \rangle = \langle Uy, Ux_0 \rangle = \langle y, x_0 \rangle = 0$.

Corollary 2. *For any invariant subspace X , the unitary operator U can*

be decomposed as $U = \text{diag}(V, W)$, where $V : X \rightarrow X$ and $W : X^\perp \rightarrow X^\perp$ are unitary operators.

Consider representation of an orthogonal projector P in the form used in **C1** to define the corresponding state transformer:

$$P = M^*M. \quad (8)$$

Let us fix some basis and represent all operators by matrices. It is convenient to select a basis in which P is diagonal $P = \text{diag}(1, \dots, 1, 0, \dots, 0)$. By the singular value decomposition theorem any matrix M can be represented in the form

$$M = W\Sigma V^*, \quad (9)$$

where the matrices W and V are unitary and Σ is a diagonal matrix such that its elements are square roots of the eigenvalues of the matrix M^*M . Hence, in the case $M^*M = P$ the matrix Σ coincides with P . Thus, in the representation (8) we can always select

$$M = WPV^*. \quad (10)$$

Moreover, we have $P = M^*M = VPV^*$ or $VP = PV$ and $V^*P = PV^*$. Thus the representation (10) can be written as

$$M = WV^*P. \quad (11)$$

We remark that the composition of two unitary operators is a unitary operator. Thus we obtained the following mathematically simple (but very useful for our further studies) result

Lemma 2. *All representations of the projector P in the form (8) are given by operators having the form:*

$$M = UP, \quad (12)$$

where U is a unitary operator.

4 Main results on combination of adjacent and separated reproducibilities with the order effect

Let $M = UP$, where P is a projector and U is a unitary operator. Set $H_P = PH$. Set $A = M^*M = P$. So, the observable A is given by a projector

(Hermitian!), but the corresponding state transformer is not of the Lüders type (not measurement of the first kind):

$$\psi \rightarrow \phi_A = \frac{UP\psi}{\|UP\psi\|}. \quad (13)$$

Proposition 1. *Observable $A = P$ with the state transformer (13) has the property of $A - A$ repeatability iff*

$$PUP = UP \quad (14)$$

Proof.

$$p_{A-A} = \langle P\phi_A, \phi_A \rangle = 1.$$

Then, as we know,

$$P\phi_A = \phi_A.$$

Thus, $PM\psi = M\psi$ or $PUP = UP$.

Corollary 2. *Observable $A = P$ with the state transformer (13) has the property of $A - A$ repeatability iff the subspace H_P is invariant with respect to the action of the unitary operator U .*

Now we consider two observables:

Let $M = U_1 P_1$ and $N = U_2 P_2$, where $P_j, j = 1, 2$, are projectors and U_j are unitary operators. Set $H_j = P_j H$ and $H_{12} = H_1 \cap H_2$.

Set $A = M^*M = P_1 U_1^* U_1 P_1 = P_1$ and $B = N^*N = P_2 U_2^* U_2 P_2 = P_2$. So, the observables A and B are given by projectors P_j (Hermitian!), but the corresponding state transformers are not of the Lüders type (these are not measurements of the first kind):

Proposition 2. *A pair of observables A and B exhibits the order effect iff*

$$P_1 U_1^* P_2 U_1 P_1 \neq P_2 U_2^* P_1 U_2 P_2 \quad (15)$$

Proof. For the $A - B$ measurement sequence, we have:

$$p_{A-B} = \|\phi_A\|^2 \langle P_2 \phi_A, \phi_A \rangle = \langle P_2 U_1 P_1 \psi, U_1 P_1 \psi \rangle = \langle P_1 U_1^* P_2 U_1 P_1 \psi, \psi \rangle;$$

in the same way, for the $B - A$ measurement sequence, we obtain:

$$p_{B-A} = \langle P_2 U_2^* P_1 U_2 P_2 \psi, \psi \rangle;$$

hence, these sequential measurements can give different results iff (15) holds true.

Proposition 3. *A pair of observables A and B exhibits the $A - B - A$ repeatability iff*

$$P_1NM = NM, \text{ or } P_1U_2P_2U_1P_1 = U_2P_2U_1P_1. \quad (16)$$

Proof. For the $A - B - A$ measurement sequence, we have :

$$p_{A-B-A} = \langle P_1\phi_{A-B}, \phi_{A-B} \rangle,$$

where

$$\phi_{A-B} = \frac{NM\psi}{\|NM\psi\|}.$$

The condition $p_{A-B-A} = 1$ is equivalent to the condition $P_1\phi_{A-B} = \phi_{A-B}$, or (16).

Thus, we are looking for observables satisfying conditions (14)–(16).

For simplicity, set $U_2 = I$, i.e., B is measurement of the first kind. Then we have the system of relations

$$P_1U_1P_1 = U_1P_1 \text{ or equivalently } U_1 : H_1 \rightarrow H_1; \quad (17)$$

$$P_1U_1^*P_2U_1P_1 \neq P_2P_1P_2; \quad (18)$$

$$P_1P_2U_1P_1 = P_2U_1P_1. \quad (19)$$

We shall show by an example that these conditions can be jointly satisfied. This statement can be reformulated in the form of the following proposition playing an important role in applications to cognition [8]–[19].

Proposition 4. *There exist quantum observables A and B generating the order effect and satisfying the conditions of $A - A$, $B - B$, and $A - B - A$ repeatability.*

Example. Consider four dimensional Hilbert space with some orthonormal basis (e_1, e_2, e_3, e_4) . Let H_1 and H_2 have bases (e_1, e_2, e_3) and (e_1, e_2, e_4) , respectively, and let P_1 and P_2 be projectors on these subspaces. It is crucial that $H_{12} = H_1 \cap H_2$ (with the basis (e_1, e_2)) is nontrivial and at least two dimensional.

Let $U = (u_{ij})$ be a unitary operator acting in H_1 . It is extended on H as $U_1 = \text{diag}(U, I)$, where I is the unit operator in H_1^\perp and (e_4) is the basis in the latter space. Thus the condition (17) is satisfied.

We check now condition (19). For any $\psi \in H$, the vector $U_1P_1\psi \in H_1$. But $P_2 : H_1 \rightarrow H_{12}$. Hence, $P_2U_1P_1\psi \in H_{12}$ and action onto this vector by P_1 cannot change it.

Finally, we show that the condition (18) holds. The crucial point is that the right-hand side does not contain the unitary operator U_1 , but the left-hand side contains it. The tricky point is that the left-hand side contains both U_1 and its inverse U_1^* . Thus, in principle, they can compensate actions of each other (cf. with the proof of Proposition 7). We shall see that this is not the case.

Let $\psi = \sum_{j=1}^4 x_j e_j$. Then $\phi = P_2 P_1 P_2 \psi \in H_{12}$ and $\phi = x_1 e_1 + x_2 e_2$.

For the other side of (18), we have: $U_1 P_1 \psi = \sum_{j=1}^3 y_j e_j$, where $y_j = \sum_{i=1}^3 u_{jix_i}$. Now consider the next step: $P_2 U_1 P_1 \psi = y_1 e_1 + y_2 e_2$. Then

$$\phi' = U_1^* P_2 U_1 P_1 \psi = \sum_{k=1}^3 z_k e_k,$$

where $z_k = \sum_{j=1}^2 \bar{u}_{jk} y_j = \sum_{j=1}^2 \sum_{i=1}^3 \bar{u}_{jk} u_{ji} x_i = \sum_{i=1}^3 (\sum_{j=1}^2 \bar{u}_{jk} u_{ji}) x_i$. Finally, since $\phi' \in H_1$, it cannot be changed by the action of P_1 .

We remark that, for a unitary matrix U , $\sum_{j=1}^3 \bar{u}_{jk} u_{ji} = \delta_{ki}$. If the subspace H_{12} is also invariant for U_1 (and, hence, for its block U), i.e., e.g., $U = \text{diag}(W, 1)$, where W is a unitary operator in H_{12} , then $\sum_{j=1}^2 \bar{u}_{jk} u_{ji} = \delta_{ki}$. Hence, in this case $\phi' = x_1 e_1 + x_2 e_2 = \phi$. However, if H_{12} is not invariant, then in general $\phi' \neq \phi$.

Now about the possibility to satisfy both $A - B - A$ and $B - A - B$ repeatability in combination with the order effect. As we have seen in the above example, for a nontrivial order effect it is crucial that the unitary operator does not leave the space $H_{12} = H_1 \cap H_2$ invariant. Now we shall take this into account in the general case. We turn again to the general case of two unitary operators U_1, U_2 . We collect the list of conditions:

1. $P_1 U_1^* P_2 U_1 P_1 \neq P_2 U_2^* P_1 U_2 P_2$;
2. $P_1 U_2 P_2 U_1 P_1 = U_2 P_2 U_1 P_1$;
3. $P_2 U_1 P_1 U_2 P_2 = U_1 P_1 U_2 P_2$.

Proposition 5. *The combination of $A - A, B - B$, and $A - B - A$ conditions of sequential repeatability implies that*

$$P_2 H_1 = H_{12} \tag{20}$$

and

$$U_2 H_{12} = H_{12}. \tag{21}$$

Proof. We start with remarking that $P_1H = H_1$ and $U_1P_1H = H_1$ as well (Corollary 2 for the condition $A - A$). Then, we have that $P_2U_1P_1H = P_2H_1 \subset H_2$ and, hence, $U_2P_2U_1P_1H = U_2(P_2H_1) \subset H_2$ (Corollary 2 for the condition $B - B$). However, the second equality in the above list of conditions implies that $U_2P_2U_1P_1H = U_2(P_2H_1) \subset H_1$ as well. Thus, $U_2(P_2H_1) \subset H_1 \cap H_2 = H_{12}$. Now we play by using finite dimensionality of the model. We remark that $H_{12} \subset H_1$ and, hence, $H_{12} \subset P_2H_1$. We remark that the dimensions of the subspaces P_2H_1 and $U_2P_2H_1$ coincide. Since $H_{12} \subset P_2H_1$ and $U_2(P_2H_1) \subset H_{12}$, We obtain that $P_2H_1 = H_{12}$ and $U_2 : H_{12} \rightarrow H_{12}$, i.e., H_{12} is invariant subspace of the unitary operator U_2 .

Lemma 3. *The condition (20) is equivalent to commutativity of projectors:*

$$[P_1, P_2] = 0. \quad (22)$$

Proof. a). Let (20) hold. Then, for any $\psi \in H$, we have $P_2P_1\psi \in H_{12}$. Hence, $P_1P_2P_1\psi = P_2P_1\psi$. Thus,

$$P_1P_2P_1 = P_2P_1. \quad (23)$$

Apply the operation of adjoint operator to the both sides of this equality. We get:

$$P_1P_2P_1 = P_1P_2. \quad (24)$$

Hence, we derived (22).

b). Let (22) hold. Then $P_2H_1 = P_2P_1H = P_1P_2H \subset H_1$ and, hence, $P_2H_1 \subset H_{12}$.

In the same way the $B - A - B$ condition implies that

$$P_1H_2 = H_{12} \quad (25)$$

and

$$U_1H_{12} = H_{12}. \quad (26)$$

Condition (25) is as well equivalent to commutativity of projectors. In particular, the conditions (20) and (25) are equivalent. The above results can be presented in the form of Proposition useful for our cognitive applications, see, e.g., [8]-[19].

Proposition 6. *The combination of the sequential repeatability conditions $A - A$ and $B - B$ with $A - B - A$ or $B - A - B$ implies that the projectors commute.*

We emphasize that $A - B - A$ implies only U_2 invariance of H_{12} and not U_1 invariance (as can be seen from the Example). In the same way

$B - A - B$ implies only U_1 invariance of H_{12} and not U_2 invariance. Thus, commutativity of projectors or conditions (20) or (25) do not imply $A - B - A$ (nor $B - A - B$) repeatability.

We can represent $H_j = H_{12} \oplus L_j$, $j = 1, 2$, where L_j is the orthogonal complement of H_{12} in the subspace H_j . In other terms $L_j = P_j H \cap (P_1 H \cap P_2 H)^\perp$. We also can decompose total Hilbert space:

$$H = H_{12} \oplus L_j \oplus H_j^\perp. \quad (27)$$

This decomposition will play a crucial role in our further considerations. We now show that the sequential repeatability conditions set a rigid constraint onto the geometric inter-relation between subspaces L_1 and L_2 .

Proposition 6a. *The combination of the sequential repeatability conditions $A - A$ and $B - B$ with $A - B - A$ or $B - A - B$ implies that $L_1 \perp L_2$.*

Proof. By Proposition 5 $A - B - A$ sequential repeatability implies that $P_2 H_1 = H_{12}$. We show that, for any pair of vectors $v \in L_1$ and $w \in L_2$, they are orthogonal. We have $\langle v, w \rangle = \langle v, P_2 w \rangle = \langle P_2 v, w \rangle = 0$, because $P_2 v \in H_{12}$.

We remark that in Hilbert space geometry, subspaces H_1 and H_2 satisfying the condition $L_1 \perp L_2$ are called *perpendicular*. It is well known that P_1, P_2 are the projectors onto perpendicular subspaces iff they commute.

By Lemma 1 the unitary operators can be represented as block-diagonal operators

$$U_j = \text{diag}(V_j, W_j, T_j), \quad (28)$$

where the diagonal blocks are unitary operators acting in spaces H_{12} , L_j , and H_j^\perp , respectively. This invariance decomposition is essential for our further considerations.

Proposition 7. *The combination of the sequential repeatability conditions $A - A$ and $B - B$ with $A - B - A$ and $B - A - B$ is incompatible with the order effect.*

Proof. By Proposition 6a the Hilbert space H can be represented as the direct sum:

$$H = H_{12} \oplus L_1 \oplus L_2 \oplus \tilde{H}, \quad \tilde{H} = (H_{12} \oplus L_1 \oplus L_2)^\perp. \quad (29)$$

Take any $\psi \in H$; it can be represented as $\psi = \psi_{12} + \phi_1 + \phi_2 + \tilde{\psi}$ with components belonging the corresponding subspaces in (29). First, we calculate $P_1 U_1^* P_2 U_1 P_1 \psi$. We have $P_1 \psi = \psi_{12} + \phi_1$. Then, $U_1 P_1 \psi = V_1 \psi_{12} + W_1 \phi_1$, see (28). It is crucial that $V_1 : H_{12} \rightarrow H_{12}$, $W_1 : L_1 \rightarrow L_1$. Thus, $W_1 \phi_1 \perp L_2$. Hence, $P_2 U_1 P_1 \psi = V_1 \psi_{12}$. Then we get $U_1^* P_2 U_1 P_1 \psi = V_1^* V_1 \psi_{12} = \psi_{12}$.

In the same way we obtain that $P_2 U_2^* P_1 U_2 P_2 \psi = \psi_{12}$, i.e., no order effect.

5 Appendix: infinite-dimensional case

The following result shows that the considerations of this paper cannot be repeated in the infinite-dimensional case:

Proposition 8. *Let $E = M^*M$. If the dimension of H is infinite, then the equality $EM = M$ does not imply that E is an orthogonal projector.*

Proof. We construct an example of the operator E which is not an orthogonal projector, but the equality $EM = M$ holds. Consider an orthonormal basis $(e_j, j = 1, 2, \dots)$ in H . Take an arbitrary complex number a . We define the operator M as

$$Me_1 = ae_2, \quad Me_n = e_{n+1}, \quad n \neq 1. \quad (30)$$

It is easy to find the adjoint operator:

$$M^*e_1 = 0, \quad M^*e_2 = \bar{a}e_1, \quad M^*e_{n+1} = e_n, \quad n \neq 1. \quad (31)$$

Then we have

$$Ee_1 = M^*Me_1 = aM^*e_2 = |a|^2e_1, \quad Ee_n = M^*Me_n = M^*e_{n+1} = e_n, \quad n \neq 1, \quad (32)$$

and

$$EMe_1 = aEe_2 = ae_2 = Me_1, \quad EMe_n = Ee_{n+1} = e_{n+1} = Me_n, \quad n \neq 1 \quad (33)$$

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